

# Finite Element Modeling of Underwater Acoustic Environments

Based on my MS Thesis

General Ozochiawaeze

Department of Mathematics  
Purdue University  
Student Colloquium

November 30, 2022

# Overview

- 1 Research Motivation
- 2 Numerical Examples
- 3 Domain Decomposition Approach
- 4 Conclusion and Future Work

# Forward Modeling

**Forward problem:** determine a model for acoustic propagation in an underwater environment consisting of a water layer and sediment layers.<sup>1</sup>

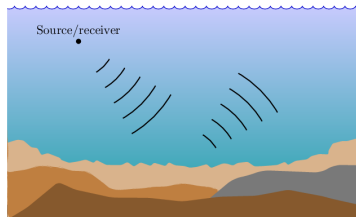
Models of high frequency  
underwater acoustics

Geometric optics

Helmholtz equations

Helmholtz equations

with attenuation



<sup>1</sup>Engquist et al., “Seafloor identification in sonar imagery via simulations of Helmholtz equations and discrete optimization”.

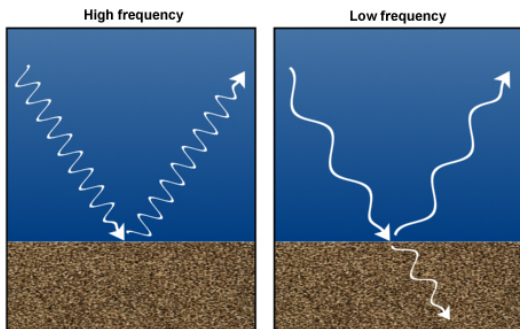
# Motivation

- This work incorporates finite element modeling of high frequency wave propagation in underwater acoustic environments.
- The forward modeling of underwater acoustic environments has important applications including:
  - ① sonar imaging for target detection
  - ② seafloor mapping
  - ③ sediment composition
  - ④ classification of marine habitats
  - ⑤ geoacoustic inversion
  - ⑥ noise prediction
- Want to simulate underwater acoustic wave scattering

# High Frequency vs Low Frequency

High frequency sound reflects more efficiently than low frequency sound.

## Effect of Sound Frequency on Reflectivity/Absorption



©The COMET Program

# The Helmholtz Equation

Most problems of acoustic scattering can be formulated in the frequency domain as linear elliptic boundary value problems (BVPs).

The main character: **Helmholtz equation**

$$\Delta u + k^2 u = 0$$

$u$  : acoustic pressure     $k := \frac{\omega}{c}$  : wave number     $c$  : sound speed  
 $\omega$  : frequency

# Why Helmholtz Equation?

$$\Delta u + k^2 u = 0 \quad \text{in } \Omega \subset \mathbb{R}^2$$

Why is it interesting?

# Why Helmholtz Equation?

$$\Delta u + k^2 u = 0 \quad \text{in } \Omega \subset \mathbb{R}^2$$

Why is it interesting?

- 1 Plenty of applications in computational ocean acoustics:
  - The Helmholtz equation is used, both analytically and numerically, to model underwater sound propagation.



# Why Helmholtz Equation?

$$\Delta u + k^2 u = 0 \quad \text{in } \Omega \subset \mathbb{R}^2$$

Why is it interesting?

- 1 Plenty of applications in computational ocean acoustics:
  - The Helmholtz equation is used, both analytically and numerically, to model underwater sound propagation.
- 2 difficult to solve numerically (for  $k \gg 1$ ):
  - oscillating solutions  $\implies$  expensive to approximate
  - hard to formulate solvers for high-frequency solutions

# Two Classes of Boundary Value Problems for the Helmholtz Equation

# Two Classes of Boundary Value Problems for the Helmholtz Equation

## 1 Interior Problems

- Helmholtz equation is set in a **bounded domain**  $\Omega \subset \mathbb{R}^d$ ,  $d = 2, 3$ , with Dirichlet, Neumann, or Impedance boundary conditions

# Two Classes of Boundary Value Problems for the Helmholtz Equation

## 1 Interior Problems

- Helmholtz equation is set in a **bounded domain**  $\Omega \subset \mathbb{R}^d$ ,  $d = 2, 3$ , with Dirichlet, Neumann, or Impedance boundary conditions

## 2 Exterior Problems

- Helmholtz equation is set in an **unbounded domain**.
- Most common and important example is the **scattering problem**.
- Scattering problems concern propagation of waves that collide with some object.
- Two types of acoustic scattering problems: **sound-soft** and **sound-hard**.

This work investigates both interior problems and exterior (scattering) problems, with greater emphasis on the latter.

# Acoustic Scattering

## Acoustic Scattering Boundary Value Problem

$$\left\{ \begin{array}{ll} \Delta u + k^2 u = 0 & \text{in } \mathbb{R}^2 \setminus \overline{\Omega}, \\ \lim_{|x| \rightarrow \infty} |x| \left( \partial_{|x|} u - iku \right) = 0 & \text{(SRC)} \\ u = -u_{inc} & \text{on } \partial\Omega \text{ (sound-soft),} \\ \partial_n u = -\partial_n u_{inc} & \text{on } \partial\Omega \text{ (sound-hard).} \end{array} \right.$$

$k$  : wavenumber

$\Omega$  : the scatterer

$u_{inc}(x) = -\exp(ik\hat{\theta} \cdot x)$ : the plane incident wave

$\hat{\theta}$ : incidence angle

$i = \sqrt{-1}$

SRC: Sommerfield Radiation Condition at infinity (outgoing scattered wave)

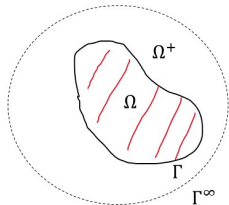
# Absorbing Boundary Conditions

## Approximation of Exterior (Scattering) Problems by Truncation

- for a numerical solution, need to truncate the unbounded domain and introduce an artificial surface with a boundary condition.
- An absorbing boundary condition (ABC) is an approximation to the Sommerfeld radiation condition.

First Order ABC:  $\frac{\partial u}{\partial n} - i\omega u = 0$  on  $\Gamma^\infty$

- $\Gamma^\infty$ : an artificial boundary.
- $\Omega^+$ : bounded domain enclosed by  $\Gamma^\infty$  and  $\Gamma$ .



# Variational Formulations

- BVPs for linear elliptic PDEs like the Helmholtz equation are often posed in variational/weak form:

$$(VF) \text{ find } u \in \mathcal{V} \text{ such that } a(u, v) = L(v) \quad \forall v \in \mathcal{V}.$$

- $\mathcal{V}$ : Hilbert space
- $a(\cdot, \cdot) : \mathcal{V} \times \mathcal{V} \rightarrow \mathbb{R}$ : bilinear form
- $L : \mathcal{V} \rightarrow \mathbb{R}$ : continuous linear functional

BVPs can be approximated using a weak Galerkin discretization:

$$(WGD) \text{ find } u_N \in \mathcal{V}_N \text{ such that } a(u_N, v_N) = L(v_N) \quad \forall v_N \in \mathcal{V}_N.$$

- $\mathcal{V}_N \subset \mathcal{V}$ : finite dimensional subspace,  $\dim(\mathcal{V}_N) = N$

# Well-Posedness of Variational Form

The weak problem is well-posed if:  $\exists c_1, c_2 > 0$  :

$$|a(u, v)| \leq c_1 \|u\|_{\mathcal{V}} \|v\|_{\mathcal{V}} \quad \forall u, v \in \mathcal{V}, \quad (\text{continuity})$$

$$|a(v, v)| \geq c_2 \|v\|_{\mathcal{V}}^2 \quad \forall v \in \mathcal{V}, \quad (\text{coercivity})$$

By the Lax-Milgram Theorem,

- well-posedness of (VF):  $\exists! u \in \mathcal{V}, \|u\|_{\mathcal{V}} \leq \frac{\|L\|_{\mathcal{V}'}}{c_1}$ ;
- well-posedness of any (WGD):  $\exists! u_N \in \mathcal{V}_N, \|u_N\|_{\mathcal{V}} \leq \frac{\|L\|_{\mathcal{V}'}}{c_1}$ ,

where  $\mathcal{V}'$  is the dual space of  $\mathcal{V}$ . Coercivity is a property of the bilinear form  $a$ .



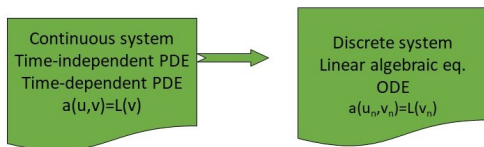
# Overview of Finite Element Method

Finite element method (FEM) is a useful numerical method for modeling underwater acoustic scattering.

Main idea: take a geometrically complex domain  $\Omega$  and divide it into cells and use polynomials (e.g., Lagrange or Hermite polynomials) for approximating a function over a cell.

## Implementation:

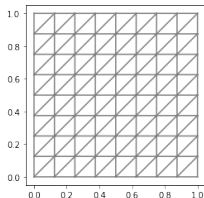
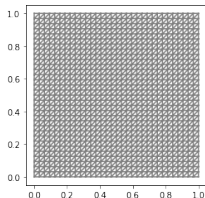
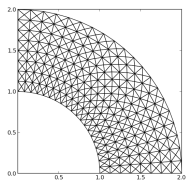
- 1 FEnICS computing platform
- 2 COMSOL Multiphysics Acoustics Module



# Overview of Finite Element Method

A finite element is defined as a triple  $(K, \mathcal{P}, \mathcal{N})$ :

- $K \subseteq \mathbb{R}^n$ : the **element domain**, a closed, bounded set with nonempty interior and a piecewise smooth boundary
- $\mathcal{P} \subset C(K)$ : the **shape functions**, a finite-dimensional space of continuous functions on  $K$
- $\mathcal{N}$ : set of **nodal variables**, an indexed family of linear functionals on  $\mathcal{P}$



**Figure:** examples of piecewise linear 2D elements on a mesh

# Numerical Results for Interior Problems

Consider the Dirichlet Helmholtz BVP

$$\begin{cases} \Delta u + k^2 u = -f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

Weak formulation is

$$\begin{cases} a(u, v) := \int_{\Omega} (\nabla u \cdot \nabla v - k^2 uv) \, dx, \\ L(v) := \int_{\Omega} f v \, dx, \\ \mathcal{V} := H_0^1(\Omega) = \{v \in H^1(\Omega) : v|_{\partial\Omega} = 0\}. \end{cases}$$

# How was the Helmholtz Weak Form Obtained?

(VF) was obtained by

- 1 multiplying  $\mathcal{L}u := \Delta u + k^2 u = f$  with **test function**  $v$ ;
- 2 using **Green's 1st identity**:

$$\int_{\Omega} [v\Delta u + \nabla v \cdot \nabla u] dx = \int_{\partial\Omega} v\nabla u \cdot n dS$$

- 3 integrating by parts
- 4 substituting the boundary condition in the boundary term

# Is the Helmholtz coercive?

- Continuity of  $a(u, v)$  follows from the Cauchy-Schwarz inequality.

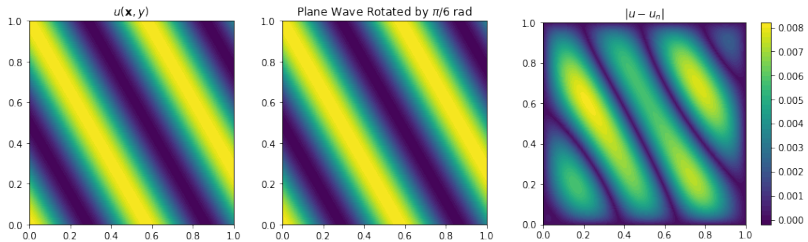
# Is the Helmholtz coercive?

- Continuity of  $a(u, v)$  follows from the Cauchy-Schwarz inequality.
- For  $k^2 \geq \lambda_1 > 0$  (1st Laplace-Dirichlet eigenvalue),  $a(\cdot, \cdot)$  is continuous but **not coercive** in  $H_0^1(\Omega)$ .
- Helmholtz equation is not coercive for large  $k$ . This makes high frequency problems hard!
- The Helmholtz problem is well-posed if  $k^2$  is not an eigenvalue of  $-\Delta$ .

# Numerical Results for Interior Problems

## Dirichlet Problem:

FEM Solution    True Solution    Plot of Absolute Error



# Numerical Results for Interior Problems

Consider also the Neumann Helmholtz BVP

$$\begin{cases} \Delta u + k^2 u = -f & \text{in } \Omega, \\ \nabla u \cdot \mathbf{n} = 0 & \text{on } \partial\Omega. \end{cases}$$

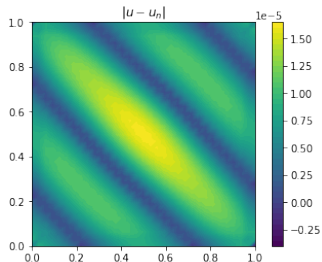
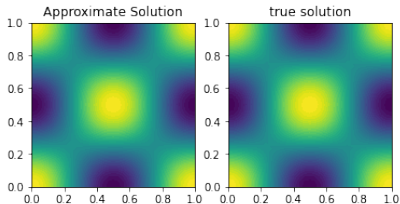
Weak formulation is

$$\begin{cases} a(u, v) := \int_{\Omega} (\nabla u \cdot \nabla v - k^2 uv) \, dx, \\ L(v) := \int_{\Omega} f v \, dx, \\ \mathcal{V} := H^1(\Omega) = \{v \in L^2(\Omega) : \partial_{x_i} v \in L^2(\Omega), 1 \leq i \leq 2\}. \end{cases}$$



# Numerical Results for Interior Problems

## Neumann Problem:



# Numerical Results for Interior Problems

**Table:** Error at the Vertex Values of Mesh N with  $P1$  elements

<b>N</b>	$\frac{\ u-u_n\ _{L^2}}{\ u\ _{L^2}}$	$\frac{\ u-u_n\ _{\infty}}{\ u_n\ _{\infty}}$
32	0.005409	0.00416
64	0.001358	0.00103
128	0.000339	0.000259
512	0.000021253	0.000001623
1024	$5.313 \times 10^{-6}$	$4.058 \times 10^{-6}$

**Table:** Error at the Vertex Values of Mesh N with  $P2$  elements

<b>N</b>	$\frac{\ u-u_n\ _{L^2}}{\ u\ _{L^2}}$ :	$\frac{\ u-u_n\ _{\infty}}{\ u_n\ _{\infty}}$ :
32	$6.93 \times 10^{-5}$	$1.28 \times 10^{-5}$
64	$8.62 \times 10^{-6}$	$8.36 \times 10^{-7}$
128	$1.07 \times 10^{-6}$	$5.29 \times 10^{-8}$
512	$1.68 \times 10^{-8}$	$2.08 \times 10^{-10}$
1024	$4.41 \times 10^{-9}$	$3.73 \times 10^{-11}$

# Point Source Problem: Helmholtz Equation With Attenuation Term

The next interior problem that serves as a model problem for generating circular waves is a Dirichlet Helmholtz BVP with a point source in the domain and zero boundary conditions on the square:

$$\begin{cases} \Delta u + (k^2 + ik\alpha)u = -\delta(x - y), \\ u = 0, \end{cases}$$

where  $ik\alpha$  denotes the attenuation term and  $\alpha$  is the attenuation/damping coefficient. The subsequent model produces waves that begin to "die out" before reaching the boundary.

# Weak Formulation of Point Source Problem

For implementation, we split the solution  $u = u_r + iu_i$ , where  $\Re u = u_r$ ,  $\Im u = u_i$ .

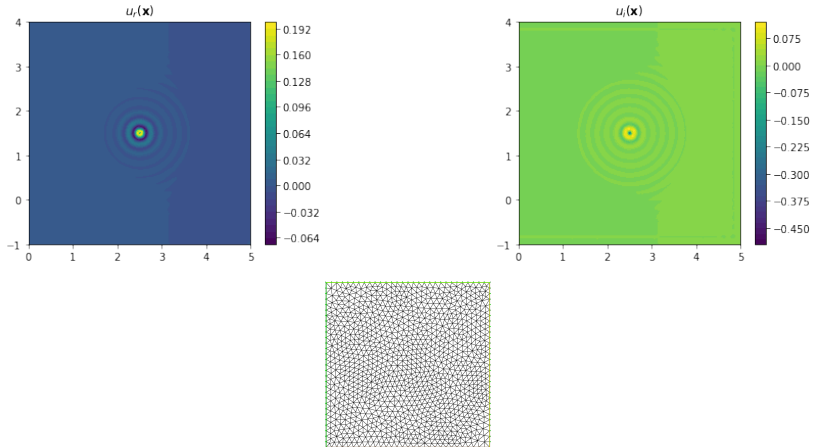
(VF) of Point Source Problem:

$$\begin{cases} a(u, v) := \int_{\Omega} (\nabla u_r \cdot \nabla v_r) + k^2 u_r v_r - \alpha k u_i v_r + (\nabla u_i \cdot \nabla v_i) + k^2 u_i v_i + \\ \alpha k u_r v_i \, d\mathbf{x}, \\ L(u, v) := \int_{\Omega} f v_r \, d\mathbf{x}, \\ \mathcal{V} := H_0^1(\Omega) \times H_0^1(\Omega) = \{(v_r, v_i) \in H^1(\Omega) : v_r|_{\partial\Omega} = v_i|_{\partial\Omega} = 0\}, \end{cases}$$

with  $u = u_r + iu_i$  and  $v = v_r + iv_i \in \mathbb{C}$ . The standard Lagrange family of elements is used via a continuous Galerkin method.

# Plot of Solution to Point Source Problem

Figure: Left:  $\Re(u)$ ; Middle: Mesh ( $N = 100$ ); Right:  $\Im(u)$

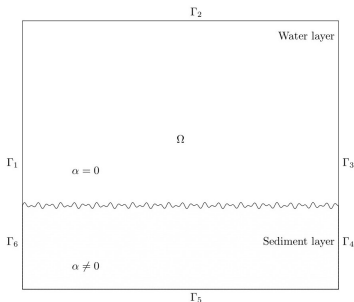


# Damped Helmholtz equation $\Delta u + \omega^2 u + i\omega\alpha u = 0$

$$u = 0 \text{ on } \Gamma_5 \quad (\text{Dirichlet})$$

$$\frac{\partial u}{\partial \mathbf{n}} - i\omega u = i\omega(1 - \sin \theta)u_{inc} \text{ on } \Gamma_2 \quad (\text{ABC})$$

$$u|_{\Gamma_1} = u|_{\Gamma_3}, \quad u|_{\Gamma_6} = u|_{\Gamma_4}$$



- $u = u_{scat} + u_{inc}$
- $u_{scat}$ : the scattered wave solution
- $u_{inc} = \exp(i\omega(\cos \theta x + \sin \theta y))$ : the incoming plane wave

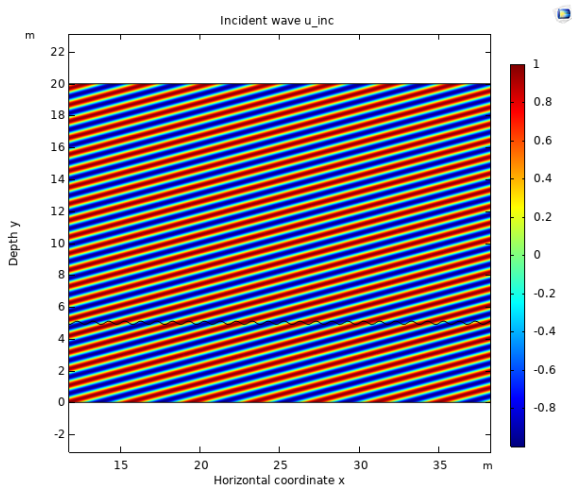
# Weak Form of Water-Sediment Model

The weak formulation of the water-sediment model is

$$\begin{cases} a(u, v) := \int_{\Omega} (-\nabla u \cdot \overline{\nabla v} + \omega^2 u \bar{v} + i\omega\alpha u) dx, \\ L(u, v) := \int_{\Gamma_2} h \bar{v} ds, \\ \mathcal{V} := H^1(\Omega), \end{cases}$$

where  $h = \nabla u_{inc} \cdot \mathbf{n}$ .

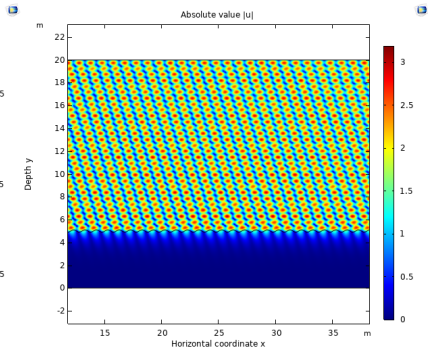
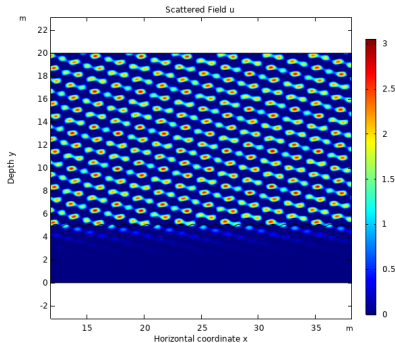
# Plot of Incident Plane Wave





# Simulation of Water-Sediment Domain

Sound attenuation occurs in the sediment.



# Why Domain Decomposition?

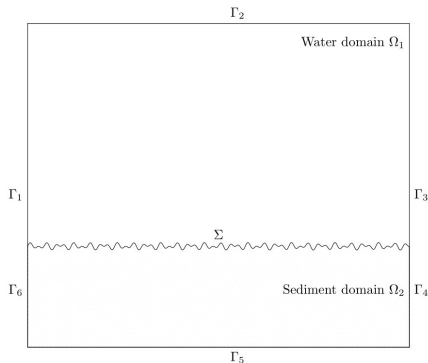
## Main difficulties

- Helmholtz equation: highly-indefinite (not coercive!) at high wave numbers
- FEM leads to large, complex-valued, and highly indefinite sparse matrix for high wave numbers
- At high wave numbers/frequencies: solution will show non-robust behavior (pollution effect). Solutions become highly oscillatory at higher wave numbers

⇒ **Alternative:** Domain Decomposition Methods!

# Lions-Desprès DDM

**Lions-Desprès DDM:** a non-overlapping domain decomposition method used to solve high frequency acoustic scattering problems (Boubendir et al.). We split the domain of the water-sediment model into two non-overlapping subdomains  $\Omega = \Omega_1 \cup \Omega_2$ .



# Lions-Deprès DDM Applied to the Water-Sediment Model

$$\left\{ \begin{array}{l} \Delta u_1 + \omega^2 u_1 = 0 \quad (\text{on } \Omega_1) \\ \Delta u_2 + (\omega^2 + i\alpha\omega) u_2 = 0 \quad (\text{on } \Omega_2) \\ \text{for the } n\text{th-iteration: } u_1, u_2 \text{ solve (SP) with the following condition on } \Sigma: \\ \partial_{n_1} u_1^n - i\omega u_1 = g_{12}^n \quad \text{on } \Sigma \\ \partial_{n_2} u_2^n - i\omega u_2 = g_{21}^n \quad \text{on } \Sigma \end{array} \right.$$

## Weak Formulation of Water Subdomain

At the  $n$ th-iteration, we multiplied by a test function  $v \in \mathcal{V}$  and found  $u_1 \in \mathcal{V}$  such that

$$\begin{cases} \mathcal{V} = \left\{ v \in H^1(\Omega) : v|_{\Gamma_1} = v|_{\Gamma_3} \right\}, \\ a(u_1, v) = \int_{\Omega_1} \nabla u_1 \cdot \nabla v - \omega^2 u_1 v \, dx - i\omega \int_{\Sigma} u_1 v \, dS, \\ L(v) = \int_{\Gamma_2} h v \, dS + \int_{\Sigma} g_{12} v \, dS, \end{cases}$$

where  $h = \nabla u_{inc} \cdot \mathbf{n}$ .

## Weak Formulation of the Sediment Subdomain

At the  $n$ th-iteration, we multiplied by a test function  $v \in \mathcal{W}$  and found  $u_2 \in \mathcal{W}$  such that

$$\left\{ \begin{array}{l} \mathcal{W} = \left\{ v \in H^1(\Omega) : v|_{\Gamma_6} = v|_{\Gamma_4} \text{ and } v|_{\Gamma_5} = 0 \right\}, \\ a(u_2, v) = \int_{\Omega_2} \nabla u_2 \cdot \nabla v - (\omega^2 + i\omega\alpha)u_2 v \, dx - i\omega \int_{\Sigma} u_2 v \, dS \\ L(v) = \int_{\Sigma} g_{21} v \, dS. \end{array} \right.$$

# Numerical Result with Lions-Despres

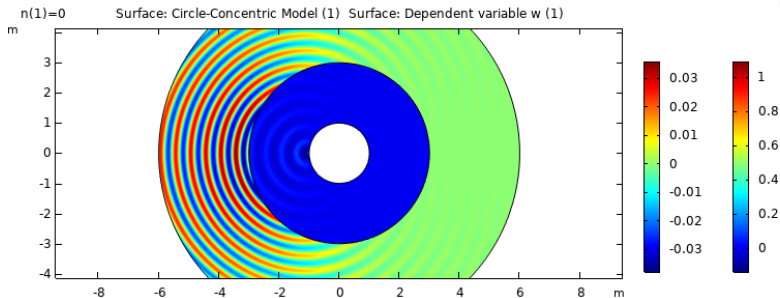


Figure: A 2d circle-concentric model of a plane wave by a unit sound-soft circular cylinder.

# Conclusion and Future Work

## Conclusion

- Forward modeling of both interior and exterior BVPs for the Helmholtz equation, simulating plane and circular waves arising in sonar imaging and noise detection
- Simulated a two-layer seafloor with attenuation in the sediment
- Formulation of the forward problem with Lions-Deprès DDM, which paves the way for detailed simulations of seafloor scattering with high frequencies



# Conclusion and Future Work

## Future Work

- implementation
- investigate the impact of the domain decomposition method on computational efficiency
- what is the number of iterations necessary to improve detailed simulation of underwater acoustic wave scattering?

# References



Yassine Boubendir, Xavier Antoine, Christopher Geuzaine (2012)

A quasi-optimal non-overlapping domain decomposition algorithm for the helmholtz equation.

*Journal of Computational Physics* 231:262 – 280.



B. Engquist, C. Frederick, Q. Huynh, and H. Zhou (2017)

Seafloor Identification in sonar imagery via simulations of Helmholtz equations and discrete optimization.

*Journal of Computational Physics* 338:477-492, June 2017.



Paul C. Etter (2012)

Advances in Acoustic Sensing, Imaging, and Signal Processing.

*Advances in Acoustics and Vibration* vol. 2012, Article ID 214839, 28 pages.

<https://doi.org/10.1155/2012/214839>



Olof Runberg (2012)

Helmholtz Equation and high frequency approximations. 2012-04.

# References



Heidi Vosbein (2010)

Introduction to Ocean Acoustics

*University Corporation for Atmospheric Research (UCAR), COMET Program.*

# Thanks for Listening!